From unbalanced optimal transport to the Camassa-Holm equation

> François-Xavier Vialard

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Ceremade, Université Paris-Dauphine INRIA team Mokaplan

Brenier's 60th birthday, Institut Henri Poincaré

Brenier's 60th birthday

Talk based on:

- P1 Unbalanced Optimal Transport: Geometry and Kantorovich formulation, with L. Chizat, B. Schmitzer, G. Peyré. (2015)
- P2 From unbalanced optimal transport to the Camassa-Holm equation, with T. Gallouet. (2016)



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Arnold's remark on incompressible Euler

Sur la géométrie différentielle des groupes de Lie de dimension infinie et ses applications à l'hydrodynamique des fluides parfaits, Ann. Inst. Fourier, 1966.

Theorem

The incompressible Euler equation is the geodesic flow of the (right-invariant) L^2 Riemannian metric on SDiff(M) (volume preserving diffeomorphisms).

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Arnold's remark on incompressible Euler

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- An intrinsic point of view by Ebin and Marsden, *Groups of diffeomorphisms and the motion of an incompressible fluid*, Ann. of Math., 1970. Short time existence results for smooth initial conditions.
- An extrinsic point of view by Brenier, relaxation of the variational problem, optimal transport, polar factorization.

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Arnold's remark continued

The incompressible Euler equation on M (Eulerian form),

$$\left\{ egin{aligned} &\partial_t v(t,x) + v(t,x) \cdot
abla v(t,x) = -
abla p(t,x), \quad t > 0, \; x \in M\,, \ &\operatorname{div}(v) = 0\,, \ &v(0,x) = v_0(x)\,, \end{aligned}
ight.$$

is the Euler-Lagrange equation for the action

$$\int_0^1 \int_M |v(t,x)|^2 \,\mathrm{d}x \,\mathrm{d}t\,,\tag{2}$$

under the flow constraint

$$\partial_t \varphi(t, x) = v(t, \varphi(t, x)),$$

div $(v) = 0.$

and time boundary value constraints:

$$\varphi(0,\cdot) = \varphi_0 \in \text{SDiff}(M) \text{ and } \varphi(1,\cdot) = \varphi_1 \in \text{SDiff}(M). \tag{3}$$

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Arnold's remark continued

Rewritten in terms of the flow φ , the action reads

$$\int_0^1 \int_M |\partial_t \varphi(t,x)|^2 \,\mathrm{d}x \,\mathrm{d}t\,,$$

under the constraint

$$\varphi(t) \in \text{SDiff}(M) \text{ for all } t \in [0,1].$$
 (5)

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Arnold's remark continued

Rewritten in terms of the flow φ , the action reads

$$\int_0^1 \int_M |\partial_t \varphi(t, x)|^2 \, \mathrm{d}x \, \mathrm{d}t \,,$$

under the constraint

$$\varphi(t) \in \text{SDiff}(M) \text{ for all } t \in [0,1].$$
 (5)

Riemannian submanifold point of view:

Let $M \hookrightarrow \mathbb{R}^d$ be isometrically embedded: A smooth curve $c(t) \in M$ is a geodesic if and only if $\ddot{c} \perp T_c M$.

Incompressible Euler in Lagrangian form:

$$\begin{cases} \ddot{\varphi} = -\nabla p \circ \varphi \\ \varphi(t) \in \text{SDiff}(M) \,. \end{cases}$$
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About Brenier's approach to incompressible Euler

Variational approach to minimizing geodesics on SDiff(M) isometrically embedded in a Hilbert space.

 Projection onto SDiff(R^d) leads to his polar factorization theorem:

Polar factorization, Y. Brenier 1991

Let $\psi \in L^2(\mathbb{R}^d, \mathbb{R}^d)$ s.t. $\psi_*(\text{Leb}) \ll \text{Leb}$, then there exists a unique couple (p, φ) (up to cste on p) s.t.

$$\psi = \nabla \boldsymbol{p} \circ \varphi,$$

and $\varphi_*(\text{Leb}) = \text{Leb}$ and p is a convex function. Moreover,

$$\|\psi - \varphi\|_{L^2} = \inf_{f} \{\|\psi - f\|_{L^2} : f_*(\text{Leb}) = \text{Leb}\}$$
(8)

- Smooth solutions of Euler are minimizing (for $t \in [0, 1]$) if $\nabla^2 p$ is bounded in L^{∞} (by π).
- In general, relaxation of the boundary value problem as (infinite) multimarginal optimal transport.

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A geometric picture: Otto's Riemannian submersion



$$\pi(\varphi) = \varphi_*(\mu)$$

 $(Dens_p(M), W_2) \quad \mu$

Figure – A Riemannian submersion: SDiff(M) as a Riemannian submanifold of $L^2(M, M)$: Incompressible Euler equation on SDiff(M)

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Reminders: Riemannian submersion

Let (M, g_M) and (N, g_N) be two Riemannian manifolds and $f: M \mapsto N$ a differentiable mapping.

Definition

The map f is a Riemannian submersion if f is a submersion and for any $x \in M$, the map $df_x : \text{Ker}(df_x)^{\perp} \mapsto T_{f(x)}N$ is an isometry.

- Vert_x := Ker(df(x)) is the vertical space.
- Hor_x $\stackrel{\text{\tiny def.}}{=}$ Ker $(df(x))^{\perp}$ is the horizontal space.
- Geodesics on N can be lifted "horizontally" to geodesics on M.

Theorem (O'Neill's formula)

Let f be a Riemannian submersion and X, Y be two orthonormal vector fields on M with horizontal lifts \tilde{X} and \tilde{Y} , then

$$\mathcal{K}_{N}(X,Y) = \mathcal{K}_{M}(\tilde{X},\tilde{Y}) + \frac{3}{4} \|\operatorname{vert}([\tilde{X},\tilde{Y}])\|_{M}^{2}, \qquad (9)$$

where K denotes the sectional curvature and vert the orthogonal projection on the vertical space.

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A pre-formulation of the polar factorization

Diff(M) $L^2(M, M)$ g1 ld SDiff(M)

$$\pi(\varphi) = \varphi_*(\mu)$$

$$(\mathsf{Dens}_p(\mathcal{M}), \mathsf{W}_2) \quad \mu$$

Figure – A Riemannian submersion: SDiff(M) as a Riemannian submanifold of $L^2(M, M)$: Incompressible Euler equation on $SDiff(M) \circ \circ \circ$

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A pre-formulation of the polar factorization

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$$\pi(arphi)=arphi_*(\mu)$$

(Dens_p(M), W₂) μ $\pi(g_1) = \mu_1$

Figure – A pre polar factorization

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A pre-formulation of the polar factorization

Diff(M) $L^2(M, M)$ g1 ld SDiff(M)

$$\pi(arphi)=arphi_*(\mu)$$

 $(\mathsf{Dens}_{\rho}(M),\mathsf{W}_2) \quad \mu \qquad \pi(g_1) = \mu_1$

Figure – Polar factorization: $g_0 = \arg \min_{g \in \text{SDiff}} ||g_1 - g||_{L^2}$

From unbalanced optimal transport to the Camassa-Holm equation

Outline

Unbalanced optimal transport

- 2 An isometric embedding
- 3 Euler-Arnold-Poincaré equation





From unbalanced optimal transport to the Camassa-Holm equation

Contents



- 2 An isometric embedding
- 3 Euler-Arnold-Poincaré equation
- The Camassa-Holm equation as an incompressible Euler equation
- 5 Corresponding polar factorization

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Monge formulation (1781)

Let $\mu, \nu \in \mathcal{P}_+(M)$,

Minimize $\int_M c(x, \varphi(x)) d\mu$

among the map s.t. $\varphi_*(\mu) = \nu$.

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$$\mathsf{Minimize} \ \int_{\mathcal{M}} c(x,\varphi(x)) d\mu$$

among the map s.t. $\varphi_*(\mu) = \nu$.

- ill posed problem, the constraint may not be satisfied.
- Ithe constraint can hardly be made weakly closed.
- \rightarrow Relaxation of the Monge problem.

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Euler-Arnold-Poincaré equation

The Camassa-Holm equation as an incompressible Euler equation

Kantorovich formulation (1942)

Let $\mu, \nu \in \mathcal{P}_+(M)$, define D by

$$D(\mu,\nu) = \inf_{\gamma \in \mathcal{P}(M^2)} \left\{ \int_{M^2} c(x,y) \, \mathrm{d}\gamma(x,y) : \pi^1_* \gamma = \mu \text{ and } \pi^2_* \gamma = \nu \right\}$$

Existence result: c lower semi-continuous and bounded from below.

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- Also valid in Polish spaces.
- If c(x, y) = ¹/_p |x y|^p, D^{1/p} is the Wasserstein distance denoted by W_p.

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- Existence result: c lower semi-continuous and bounded from below.
- Also valid in Polish spaces.
- If c(x, y) = ¹/_p |x y|^p, D^{1/p} is the Wasserstein distance denoted by W_p.

Linear optimization problem and associated numerical methods. Recently introduced, entropic regularization. (C. Léonard, M. Cuturi, ...) From unbalanced optimal transport to the Camassa-Holm equation

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The Camassa-Holm equation as an incompressible Euler equation

Reminders: Dynamic formulation (Benamou-Brenier)

For geodesic costs, for instance $c(x, y) = \frac{1}{2}|x - y|^2$

$$\inf \mathcal{E}(v) = \frac{1}{2} \int_0^1 \int_M |v(x)|^2 \rho(x) \, \mathrm{d}x \, \mathrm{d}t \quad , \tag{11}$$

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Corresponding polar factorization

s.t.

$$egin{cases} \dot{
ho}+
abla\cdot(m{v}
ho)=0\
ho(0)=\mu_0 ext{ and }
ho(1)=\mu_1\,. \end{cases}$$

Reminders: Dynamic formulation (Benamou-Brenier)

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s.t.

$$\begin{cases} \dot{\rho} + \nabla \cdot (\boldsymbol{v}\rho) = 0\\ \rho(0) = \mu_0 \text{ and } \rho(1) = \mu_1. \end{cases}$$
(12)

Convex reformulation: Change of variable: momentum $m = \rho v$,

$$\inf \mathcal{E}(m) = \frac{1}{2} \int_0^1 \int_M \frac{|m(x)|^2}{\rho(x)} \, \mathrm{d}x \, \mathrm{d}t \,, \tag{13}$$

s.t.

$$\begin{cases} \dot{\rho} + \nabla \cdot m = 0\\ \rho(0) = \mu_0 \text{ and } \rho(1) = \mu_1. \end{cases}$$
(14)

where $(\rho, m) \in \mathcal{M}([0, 1] \times M, \mathbb{R} \times \mathbb{R}^d)$.

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s.t.

$$\begin{cases} \dot{\rho} + \nabla \cdot m = 0\\ \rho(0) = \mu_0 \text{ and } \rho(1) = \mu_1. \end{cases}$$
(14)

where $(\rho, m) \in \mathcal{M}([0, 1] \times M, \mathbb{R} \times \mathbb{R}^d)$.

Existence of minimizers: Fenchel-Rockafellar. Numerics: First-order splitting algorithm: Douglas-Rachford. From unbalanced optimal transport to the Camassa-Holm equation

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Unbalanced optimal transport

An isometric embedding

Euler-Arnold-Poincaré equation

The Camassa-Holm equation as an incompressible Euler equation

Starting point and initial motivation

- Extend the Wasserstein L² distance to positive Radon measures.
- Develop associated numerical algorithms.

Possible applications: Imaging, machine learning, gradient flows, ...

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Figure - Optimal transport between bimodal densities

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Figure – Another transformation

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Bibliography before (june) 2015

Taking into account locally the change of mass:

Two directions: Static and dynamic.

- Static, Partial Optimal Transport [Figalli & Gigli, 2010]
- Static, Hanin 1992, Benamou and Brenier 2001.
- Dynamic, Numerics, Metamorphoses [Maas et al., 2015]
- Dynamic, Numerics, Growth model [Lombardi & Maitre, 2013]

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 Dynamic and static, [Piccoli & Rossi, 2013, Piccoli & Rossi, 2014] From unbalanced optimal transport to the Camassa-Holm equation

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- Dynamic and static, [Piccoli & Rossi, 2013, Piccoli & Rossi, 2014]

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No equivalent of L^2 Wasserstein distance on positive Radon measures.

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Bibliography after june 2015

More than 300 pages on the same model!

Starting point: Dynamic formulation

- Dynamic, Numerics, Imaging [Chizat et al. , 2015]
- Dynamic, Geometry and Static [Chizat et al. , 2015]
- Dynamic, Gradient flow [Kondratyev et al. , 2015]
- Dynamic, Gradient flow [Liero et al., 2015b]
- Static and more [Liero et al. , 2015a]
- Optimal transport for contact forms [Rezakhanlou, 2015]
- Static relaxation of OT, machine learning [Frogner *et al.*, 2015]

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Two possible directions

Pros and cons:

• Extend static formulation: Frogner et al.

min
$$\lambda KL(\operatorname{Proj}_*^1 \gamma, \rho_1) + \lambda KL(\operatorname{Proj}_*^2 \gamma, \rho_2)$$

+ $\int_{M^2} \gamma(x, y) d(x, y)^2 \, \mathrm{d}x \, \mathrm{d}y$ (15)

Good for numerics, but is it a distance ?

• Extend dynamic formulation: on the tangent space of a density, choose a metric on the transverse direction. Built-in metric property but does there exist a static formulation ? From unbalanced optimal transport to the Camassa-Holm equation

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An extension of Benamou-Brenier formulation

Add a source term in the constraint: (weak sense)

$$\dot{
ho} = -\nabla \cdot (
ho \mathbf{v}) + \mathbf{\alpha} \mathbf{
ho} \,,$$

where α can be understood as the growth rate.

$$WF(m,\alpha)^2 = \frac{1}{2} \int_0^1 \int_M |v(x,t)|^2 \rho(x,t) \, \mathrm{d}x \, \mathrm{d}t + \frac{\delta^2}{2} \int_0^1 \int_M \alpha(x,t)^2 \rho(x,t) \, \mathrm{d}x \, \mathrm{d}t.$$

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where δ is a length parameter.

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where δ is a length parameter.

Remark: very natural and not studied before.

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Convex reformulation

Add a source term in the constraint: (weak sense)

$$\dot{\rho} = -\nabla \cdot \boldsymbol{m} + \boldsymbol{\mu} \,.$$

The Wasserstein-Fisher-Rao metric:

$$\mathsf{WF}(m,\mu)^2 = \frac{1}{2} \int_0^1 \int_M \frac{|m(x,t)|^2}{\rho(x,t)} \, \mathrm{d}x \, \mathrm{d}t + \frac{\delta^2}{2} \int_0^1 \int_M \frac{\mu(x,t)^2}{\rho(x,t)} \, \mathrm{d}x \, \mathrm{d}t \, .$$

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$$\mathsf{WF}(m,\mu)^2 = \frac{1}{2} \int_0^1 \int_M \frac{|m(x,t)|^2}{\rho(x,t)} \, \mathrm{d}x \, \mathrm{d}t + \frac{\delta^2}{2} \int_0^1 \int_M \frac{\mu(x,t)^2}{\rho(x,t)} \, \mathrm{d}x \, \mathrm{d}t \, .$$

- Fisher-Rao metric: Hessian of the Boltzmann entropy/ Kullback-Leibler divergence and reparametrization invariant. Wasserstein metric on the space of variances in 1D.
- Convex and 1-homogeneous: convex analysis (existence and more)
- Numerics: First-order splitting algorithm: Douglas-Rachford.
- Code available at https://github.com/lchizat/optimal-transport/

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A general framework

Definition (Infinitesimal cost)

An infinitesimal cost is $f: M \times \mathbb{R} \times \mathbb{R}^d \times \mathbb{R} \to \mathbb{R}_+ \cup \{+\infty\}$ such that for all $x \in M$, $f(x, \cdot, \cdot, \cdot)$ is convex, positively 1-homogeneous, lower semicontinuous and satisfies

$$f(x, \rho, m, \mu) \begin{cases} = 0 & \text{if } (m, \mu) = (0, 0) \text{ and } \rho \ge 0 \\ > 0 & \text{if } |m| \text{ or } |\mu| > 0 \\ = +\infty & \text{if } \rho < 0 \,. \end{cases}$$

Definition (Dynamic problem)

For $(
ho,m,\mu)\in\mathcal{M}([0,1] imes M)^{1+d+1}$, let

$$J(\rho, m, \mu) \stackrel{\text{\tiny def.}}{=} \int_0^1 \int_M f(x, \frac{\mathrm{d}\rho}{\mathrm{d}\lambda}, \frac{\mathrm{d}m}{\mathrm{d}\lambda}, \frac{\mathrm{d}\mu}{\mathrm{d}\lambda}) \,\mathrm{d}\lambda(t, x) \tag{16}$$

The dynamic problem is, for $\rho_0, \rho_1 \in \mathcal{M}_+(M)$,

$$C(\rho_0, \rho_1) \stackrel{\text{\tiny def.}}{=} \inf_{(\rho, \omega, \zeta) \in \mathcal{CE}_0^1(\rho_0, \rho_1)} J(\rho, \omega, \zeta).$$
(17)

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Existence of minimizers

Proposition (Fenchel-Rockafellar)

Let B(x) be the polar set of $f(x, \cdot, \cdot, \cdot)$ for all $x \in M$ and assume it is a lower semicontinuous set-valued function. Then the minimum of (17) is attained and it holds

$$C_{\mathcal{D}}(\rho_0,\rho_1) = \sup_{\varphi \in \mathcal{K}} \int_{\mathcal{M}} \varphi(1,\cdot) \,\mathrm{d}\rho_1 - \int_{\mathcal{M}} \varphi(0,\cdot) \,\mathrm{d}\rho_0 \qquad (18)$$

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with
$$K \stackrel{\text{\tiny def.}}{=} \{ \varphi \in C^1([0,1] \times M) : (\partial_t \varphi, \nabla \varphi, \varphi) \in B(x), \, \forall (t,x) \in [0,1] \times M \}$$

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Existence of minimizers

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$$C_{D}(\rho_{0},\rho_{1}) = \sup_{\varphi \in \mathcal{K}} \int_{\mathcal{M}} \varphi(1,\cdot) \,\mathrm{d}\rho_{1} - \int_{\mathcal{M}} \varphi(0,\cdot) \,\mathrm{d}\rho_{0} \qquad (18)$$

with
$$\mathcal{K} \stackrel{\text{\tiny def.}}{=} \{ \varphi \in \mathcal{C}^1([0,1] \times \mathcal{M}) : (\partial_t \varphi, \nabla \varphi, \varphi) \in \mathcal{B}(x), \, \forall (t,x) \in [0,1] \times \mathcal{M} \}$$

$$WF(x, y, z) = \begin{cases} \frac{|y|^2 + \delta^2 z^2}{2x} & \text{if } x > 0, \\ 0 & \text{if } (x, |y|, z) = (0, 0, 0) \\ +\infty & \text{otherwise} \end{cases}$$

and the corresponding Hamilton-Jacobi equation is

$$\partial_t \varphi + \frac{1}{2} \left(|\nabla \varphi|^2 + \frac{\varphi^2}{\delta_*^2} \right) \leq 0.$$

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Figure - WFR geodesic between bimodal densities

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Figure – Geodesics between ρ_0 and ρ_1 for (1st row) Hellinger, (2nd row) W_2 , (3rd row) partial OT, (4th row) WF.

An Interpolating Distance between Optimal Transport and Fisher-Rao, L. Chizat, B. Schmitzer, G. Peyré, and F.-X. Vialard, FoCM, 2016. From unbalanced optimal transport to the Camassa-Holm equation

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From dynamic to static

Group action

Mass can be moved and changed: consider $m(t)\delta_{x(t)}$.

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From dynamic to static

Group action

Mass can be moved and changed: consider $m(t)\delta_{x(t)}$.

Infinitesimal action

$$\dot{
ho} = -
abla \cdot (\mathbf{v}
ho) + \mu \Leftrightarrow egin{cases} \dot{x}(t) = \mathbf{v}(t, \mathbf{x}(t)) \ \dot{m}(t) = \mu(t, \mathbf{x}(t)) \ \dot{m}(t) = \mu(t, \mathbf{x}(t)) \end{cases}$$

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From dynamic to static

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Infinitesimal action

$$\dot{
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ho) + \mu \Leftrightarrow egin{cases} \dot{x}(t) = m{v}(t, x(t)) \ \dot{m}(t) = \mu(t, x(t)) \end{cases}$$

A cone metric

WF²(x,m)((
$$\dot{x}, \dot{m}$$
), (\dot{x}, \dot{m})) = $\frac{1}{2}(m\dot{x}^2 + \frac{\dot{m}^2}{m})$,

Change of variable: $r^2 = m...$

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Riemannian cone

Definition

Let (M, g) be a Riemannian manifold. The cone over (M, g) is the Riemannian manifold $(M \times \mathbb{R}^*_+, r^2g + dr^2)$. From unbalanced optimal transport to the Camassa-Holm equation

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Riemannian cone

Definition

Let (M, g) be a Riemannian manifold. The cone over (M, g) is the Riemannian manifold $(M \times \mathbb{R}^*_+, r^2g + dr^2)$.



For $M = S_1(r)$, radius $r \leq 1$. One has $sin(\alpha) = r$.

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Geometry of a cone

- Change of variable: $WF^2 = \frac{1}{2}r^2g + 2 dr^2$.
- Non complete metric space: add the vertex $M \times \{0\}$.
- The distance:

$$d((x_1, m_1), (x_2, m_2))^2 = m_2 + m_1 - 2\sqrt{m_1 m_2} \cos\left(\frac{1}{2} d_M(x_1, x_2) \wedge \pi\right) . \quad (19)$$

•
$$M = \mathbb{R}$$
 then $(x, m) \mapsto \sqrt{me^{ix/2}} \in \mathbb{C}$ local isometry.

Corollary

If (M, g) has sectional curvature greater than 1, then $(M \times \mathbb{R}^*_+, mg + \frac{1}{4m} dm^2)$ has non-negative sectional curvature. For X, Y two orthornormal vector fields on M.

$$K(\tilde{X}, \tilde{Y}) = (K_g(X, Y) - 1)$$
⁽²⁰⁾

where K and K_g denote respectively the sectional curvatures of $M \times \mathbb{R}^*_{\perp}$ and M.

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Visualize geodesics for $r^2g + dr^2$



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Distance between Diracs



$$\frac{1}{4}WF(m_1\delta_{x_1},m_2\delta_{x_2})^2 = m_2 + m_1 \\ -2\sqrt{m_1m_2}\cos\left(\frac{1}{2}d_M(x_1,x_2) \wedge \pi/2\right) \,.$$

Proof: prove that an explicit geodesic is a critical point of the convex functional.

Properties: positively 1-homogeneous and convex in (m_1, m_2) .

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Generalization of Otto's Riemannian submersion Idea of a left group action:

$$\pi: \left(\mathsf{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+)\right) \times \mathsf{Dens}(M) \mapsto \mathsf{Dens}(M)$$
$$\pi\left((\varphi, \lambda), \rho\right) := \varphi_*(\lambda^2 \rho)$$

Group law:

$$(\varphi_1, \lambda_1) \cdot (\varphi_2, \lambda_2) = (\varphi_1 \circ \varphi_2, (\lambda_1 \circ \varphi_2)\lambda_2)$$
(21)

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Generalization of Otto's Riemannian submersion Idea of a left group action:

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ho)$

Group law:

$$(\varphi_1, \lambda_1) \cdot (\varphi_2, \lambda_2) = (\varphi_1 \circ \varphi_2, (\lambda_1 \circ \varphi_2)\lambda_2)$$
(21)

Theorem (P1)

Let $\rho_0 \in \text{Dens}(M)$ and $\pi_0 : \text{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+) \mapsto \text{Dens}(M)$ defined by $\pi_0(\varphi, \lambda) := \varphi_*(\lambda^2 \rho_0)$. It is a Riemannian submersion

 $(\operatorname{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+), L^2(M, M \times \mathbb{R}^*_+)) \xrightarrow{\pi_0} (\operatorname{Dens}(M), \operatorname{WF})$

(where $M \times \mathbb{R}^*_+$ is endowed with the cone metric).

O'Neill's formula: sectional curvature of (Dens(M), WF).

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Geometric consequence

The sectional curvature of Dens(M) at point ρ is:

$$\mathcal{K}(\rho)(X_1, X_2) = \int_M k(x, 1)(Z_1(x), Z_2(x))w(Z_1(x), Z_2(x))\rho(x) \,\mathrm{d}\nu(x) \\ + \frac{3}{4} \left\| [Z_1, Z_2]^V \right\|^2 \quad (22)$$

where

$$w(Z_1(x), Z_2(x)) = g(x)(Z_1(x), Z_1(x))g(x)(Z_2(x), Z_2(x)) - g(x)(Z_1(x), Z_2(x))^2$$

and $[Z_1, Z_2]^V$ denotes the vertical projection of $[Z_1, Z_2]$ at identity and $\|\cdot\|$ denotes the norm at identity.

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Geometric consequence

The sectional curvature of Dens(M) at point ρ is:

$$\begin{split} \mathcal{K}(\rho)(X_1, X_2) &= \int_M k(x, 1)(Z_1(x), Z_2(x))w(Z_1(x), Z_2(x))\rho(x) \,\mathrm{d}\nu(x) \\ &+ \frac{3}{4} \left\| [Z_1, Z_2]^V \right\|^2 \quad (22) \end{split}$$

where

$$w(Z_1(x), Z_2(x)) = g(x)(Z_1(x), Z_1(x))g(x)(Z_2(x), Z_2(x)) - g(x)(Z_1(x), Z_2(x))^2$$

and $[Z_1, Z_2]^V$ denotes the vertical projection of $[Z_1, Z_2]$ at identity and $\|\cdot\|$ denotes the norm at identity.

Corollary

Let (M,g) be a compact Riemannian manifold of sectional curvature bounded below by 1, then the sectional curvature of (Dens(M), WF) is non-negative.

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Consequences

Monge formulation

$$WF(\rho_{0},\rho_{1}) = \inf_{(\varphi,\lambda)} \left\{ \|(\varphi,\lambda) - (Id,1)\|_{L^{2}(\rho_{0})} : \varphi_{*}(\lambda^{2}\rho_{0}) = \rho_{1} \right\}$$
(23)

Under existence and smoothness of the minimizer, there exists a function $p \in C^{\infty}(M, \mathbb{R})$ such that

$$(\varphi(x),\lambda(x)) = \exp_x^{\mathcal{C}(M)} \left(\frac{1}{2}\nabla p(x), p(x)\right), \qquad (24)$$

Equivalent to Monge-Ampère equation

With $z \stackrel{\text{\tiny def.}}{=} \log(1 + p)$ one has

$$(1+|
abla z|^2)e^{2z}
ho_0=\det(Darphi)
ho_1\circarphi$$

and

$$\varphi(x) = \exp_{(x,1)}^{M} \left(\arctan\left(\frac{1}{2}|\nabla z|\right) \frac{\nabla z(x)}{|\nabla z(x)|} \right) \,.$$

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A relaxed static OT formulation

Define

sut

$$\mathcal{KL}(\gamma, \nu) = \int \frac{\mathrm{d}\gamma}{\mathrm{d}\nu} \log\left(\frac{\mathrm{d}\gamma}{\mathrm{d}\nu}\right) \,\mathrm{d}\nu + |\nu| - |\gamma|$$

Theorem (Dual formulation, P1)

$$WF^{2}(\rho_{0},\rho_{1}) = \sup_{(\phi,\psi)\in C(M)^{2}} \int_{M} \phi(x) \,\mathrm{d}\rho_{0} + \int_{M} \psi(y) \,\mathrm{d}\rho_{1}$$

opect to $\forall (x,y) \in M^{2}, \, \phi(x) \leq 1, \quad \psi(y) \leq 1$ and
 $(1 - \phi(x))(1 - \psi(y)) \geq \cos^{2}(|x - y|/2 \wedge \pi/2)$

The corresponding primal formulation

$$WF^{2}(\rho_{1},\rho_{2}) = \inf_{\gamma} KL(\operatorname{Proj}_{*}^{1} \gamma,\rho_{1}) + KL(\operatorname{Proj}_{*}^{2} \gamma,\rho_{2})$$
$$- \int_{M^{2}} \gamma(x,y) \log(\cos^{2}(d(x,y)/2 \wedge \pi/2)) \, \mathrm{d}x \, \mathrm{d}y$$

Theorem (P2)

On a Riemannian manifold (compact without boundary), the static and dynamic formulations are equal.

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New algorithm

Scaling Algorithms for Unbalanced Transport Problems, L. Chizat, G. Peyré, B. Schmitzer, F.-X. Vialard.

• Use of entropic regularization.

$$WF^{2}(\rho_{1},\rho_{2}) = \inf_{\gamma} KL(\operatorname{Proj}_{*}^{1}\gamma,\rho_{1}) + KL(\operatorname{Proj}_{*}^{2}\gamma,\rho_{2})$$
$$- \int_{M^{2}} \gamma(x,y) \log(\cos^{2}(d(x,y)/2 \wedge \pi/2)) \, \mathrm{d}x \, \mathrm{d}y + \varepsilon KL(\gamma,\mu_{0}) \, .$$

- Alternate projection algorithm (contraction for a Hilbert type metric).
- Applications to color transfer, Fréchet-Karcher mean (barycenters).
- Simulations for gradient flows.

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The Riemannian submersion for WFR



Figure – The same picture in our case: what is the corresponding equation to Euler?

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The isotropy subgroup for unbalanced optimal transport

Recall that

$$\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \lambda) \in \mathsf{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+) : \varphi_*(\lambda^2 \rho_0) = \rho_0\}$$

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The isotropy subgroup for unbalanced optimal transport

Recall that

$$\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \lambda) \in \mathsf{Diff}(M) \ltimes C^\infty(M, \mathbb{R}^*_+) : \varphi_*(\lambda^2 \rho_0) = \rho_0\}$$

 $\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \sqrt{\operatorname{Jac}(\varphi)}) \in \operatorname{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+) : \varphi \in \operatorname{Diff}(M)\}$ The vertical space is

$$\operatorname{Vert}_{(\varphi,\lambda)} = \{(\nu,\alpha) \circ (\varphi,\lambda); \operatorname{div}(\rho\nu) = 2\alpha\rho\} , \qquad (26)$$

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where $(v, \alpha) \in Vect(M) \times C^{\infty}(M, \mathbb{R})$. The horizontal space is

$$\operatorname{Hor}_{(\varphi,\lambda)} = \left\{ \left(\frac{1}{2} \nabla p, p \right) \circ (\varphi, \lambda); \ p \in C^{\infty}(M, \mathbb{R}) \right\}.$$
(27)

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Recall that

$$\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \lambda) \in \mathsf{Diff}(M) \ltimes C^\infty(M, \mathbb{R}^*_+) \, : \, \varphi_*(\lambda^2 \rho_0) = \rho_0\}$$

 $\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \sqrt{\operatorname{Jac}(\varphi)}) \in \operatorname{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+) : \varphi \in \operatorname{Diff}(M)\}$ The vertical space is

$$\operatorname{Vert}_{(\varphi,\lambda)} = \{ (\nu, \alpha) \circ (\varphi, \lambda) ; \operatorname{div}(\rho \nu) = 2\alpha \rho \} , \qquad (26)$$

where $(v, \alpha) \in \text{Vect}(M) \times C^{\infty}(M, \mathbb{R})$. The horizontal space is

$$\operatorname{Hor}_{(\varphi,\lambda)} = \left\{ \left(\frac{1}{2} \nabla p, p \right) \circ (\varphi, \lambda); \ p \in C^{\infty}(M, \mathbb{R}) \right\}.$$
(27)

The induced metric is

$$G(\mathbf{v},\operatorname{div}\mathbf{v}) = \int_{M} |\mathbf{v}|^{2} d\mu + \frac{1}{4} \int_{M} |\operatorname{div}\mathbf{v}|^{2} d\mu.$$
 (28)

The H^{div} right-invariant metric on the group of diffeomorphisms.

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Right-invariant metric on a Lie group

Definition (Right-invariant metric)

Let $g_1, g_2 \in G$ be two group elements, the distance between g_1 and g_2 can be defined by:

$$d^2(g_1,g_2) = \inf_{g(t)} \left\{ \int_0^1 \|v(t)\|_{\mathfrak{g}}^2 \, dt \, |g(0) = g_0 \, and \, g(1) = g_1
ight\}$$

where $\partial_t g(t)g(t)^{-1} = v(t) \in \mathfrak{g}$, with \mathfrak{g} the Lie algebra.

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Right-invariant metric on a Lie group

Definition (Right-invariant metric)

Let $g_1, g_2 \in G$ be two group elements, the distance between g_1 and g_2 can be defined by:

$$d^2(g_1,g_2) = \inf_{g(t)} \left\{ \int_0^1 \|v(t)\|_{\mathfrak{g}}^2 dt \, |g(0) = g_0 \text{ and } g(1) = g_1
ight\}$$

where $\partial_t g(t)g(t)^{-1} = v(t) \in \mathfrak{g}$, with \mathfrak{g} the Lie algebra.

Right-invariance means:

$$d^2(g_1g,g_2g) = d(g_1,g_2).$$

It comes from:

$$\partial_t (g(t)g_0)(g(t)g_0)^{-1} = \partial_t g(t)g_0g_0^{-1}g(t)^{-1} = \partial_t g(t)g(t)^{-1}$$
 .

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Euler-Arnold-Poincaré equation

Compute the Euler-Lagrange equation of the distance functional:

$$\frac{\partial L}{\partial g} - \frac{d}{dt} \frac{\partial L}{\partial \dot{g}} = 0$$

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Euler-Arnold-Poincaré equation

Compute the Euler-Lagrange equation of the distance functional:

$$\frac{\partial L}{\partial g} - \frac{d}{dt} \frac{\partial L}{\partial \dot{g}} = 0$$

In the case of $\int_0^1 L(g, \dot{g}) dt = \int_0^1 ||u||^2 dt$, Euler-Poincaré-Arnold equation

$$\dot{g} = u \circ g$$

 $\dot{u} + \operatorname{ad}_{u}^{*} u = 0$ (29)

where $\operatorname{ad}_{u}^{*}$ is the (metric) adjoint of $\operatorname{ad}_{u}v = [v, u]$.

Proof.

Compute variations of v(t) in terms of $u(t) = \delta g(t)g(t)^{-1}$. Find that admissible variations on g can be written as: $\delta v(t) = \dot{u} - \mathrm{ad}_v u$ for any u vanishing at 0 and 1. From unbalanced optimal transport to the Camassa-Holm equation

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Fluid dynamics examples of Euler-Arnold equations

- Incompressible Euler equation.
- Korteweg-de-Vries equation.
- Camassa-Holm equation 1981/1993. An integrable shallow water equation with peaked solitons

Consider Diff(S_1) endowed with the H^1 right-invariant metric $\|v\|_{L^2}^2 + \frac{1}{4} \|\partial_x v\|_{L^2}^2$. One has

$$\begin{cases} \partial_t u - \frac{1}{4} \partial_{txx} u \, u + 3 \partial_x u \, u - \frac{1}{2} \partial_{xx} u \, \partial_x u - \frac{1}{4} \partial_{xxx} u \, u = 0\\ \partial_t \varphi(t, x) = u(t, \varphi(t, x)). \end{cases}$$
(30)

- Model for waves in shallow water.
- Completely integrable system (bi-Hamiltonian).
- Exhibits particular solutions named as peakons. (geodesics as collective Hamiltonian).
- Blow-up of solutions which gives a model for wave breaking.

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Ebin-Marsden analytical framework

Rewrite the metric in Lagrangian coordinates φ and a tangent vector X_{φ} and realize that it is smooth...

• The right-invariant H^{div} metric:

$$G_{\varphi}(X_{\varphi}, X_{\varphi}) = \int_{M} a^{2} |X_{\varphi} \circ \varphi^{-1}|^{2} + b^{2} \operatorname{div}(X_{\varphi} \circ \varphi^{-1})^{2} d\mu. \quad (31)$$

can be written

$$G_{\varphi}(X_{\varphi}, X_{\varphi}) = \int_{M} a^{2} |X_{\varphi}|^{2} \operatorname{Jac}(\varphi) + b^{2} \left(\operatorname{Tr}(DX_{\varphi} \cdot [D\varphi]^{-1}) \right)^{2} \operatorname{Jac}(\varphi) \, \mathrm{d}\mu \, .$$

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Smooth metric on an infinite dimensional Riemannian manifold.

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can be written

$$G_{\varphi}(X_{\varphi}, X_{\varphi}) = \int_{M} a^2 |X_{\varphi}|^2 \operatorname{Jac}(\varphi) + b^2 \left(\operatorname{Tr}(DX_{\varphi} \cdot [D\varphi]^{-1})\right)^2 \operatorname{Jac}(\varphi) \operatorname{d} \mu$$

Smooth metric on an infinite dimensional Riemannian manifold. Consequences:

- Geodesic equations is a simple ODE (No need for a Riemannian connection)
- Gauss lemma on H^s for s > d/2 + 2
- Geodesics are minimizing within H^s topology.

Theorem (Consequence of Ebin and Marsden)

Local well-posedness of the geodesics for the H^{div} right-invariant metric on Diff^s(M) for s > d/2 + 2.

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Metric properties

Theorem (Michor and Mumford, 2005)

The distance on Diff(M) endowed with the right-invariant metric L^2 is degenerate; i.e. $d(\varphi_0, \varphi_1) = 0$ for every $\varphi_0, \varphi_1 \in \text{Diff}(M)$.

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The distance on Diff(M) endowed with the right-invariant metric L^2 is degenerate; i.e. $d(\varphi_0, \varphi_1) = 0$ for every $\varphi_0, \varphi_1 \in \text{Diff}(M)$.

Theorem (Michor and Mumford, 2005)

The distance on Diff(M) endowed with the right-invariant metric H^{Div} is non degenerate.

Proof.

Direct using the isometric injection.

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We have

inj : (Diff(
$$M$$
), H^{div}) $\hookrightarrow L^2(M, \mathcal{C}(M))$
 $\varphi \mapsto (\varphi, \sqrt{\text{Jac}(\varphi)}).$

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An isometric embedding

We have

$$\begin{split} \mathsf{inj} &: (\mathsf{Diff}(M), H^{\mathsf{div}}) \hookrightarrow L^2(M, \mathcal{C}(M)) \\ & \varphi \mapsto (\varphi, \sqrt{\mathsf{Jac}(\varphi)}) \,. \end{split}$$

The geodesic equations can be written in Lagrangian coordinates

$$\begin{cases} \frac{D}{Dt}\dot{\varphi} + 2\frac{\dot{\lambda}}{\lambda}\dot{\varphi} = -\nabla^{g}P\circ\varphi\\ \ddot{\lambda}r - \lambda rg(\dot{\varphi},\dot{\varphi}) = -2\lambda rP\circ\varphi. \end{cases}$$
(32)

In Eulerian coordinates,

$$\begin{cases} \dot{\mathbf{v}} + \nabla^{g}_{\mathbf{v}} \mathbf{v} + 2\mathbf{v}\alpha = -\nabla^{g} P\\ \dot{\alpha} + \langle \nabla \alpha, \mathbf{v} \rangle + \alpha^{2} - g(\mathbf{v}, \mathbf{v}) = -2P \,, \end{cases}$$
(33)

where $\alpha = \frac{\dot{\lambda}}{\lambda} \circ \varphi^{-1}$ and $v = \partial_t \varphi \circ \varphi^{-1}$.

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Consequences of the isometric embedding

$(\mathsf{Diff}(M), H^{\mathsf{div}}) \hookrightarrow L^2(M, \mathcal{C}(M))$

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Consequences of the isometric embedding

 $(\operatorname{Diff}(M), H^{\operatorname{div}}) \hookrightarrow L^2(M, \mathcal{C}(M))$

- Using Gauss-Codazzi formula, it generalizes a curvature formula by Khesin et al. obtained on $\text{Diff}(S_1)$.
- Smooth geodesics are length minimizing for a short enough time under mild conditions (generalization of Brenier's proof).
- The Camassa-Holm equation as incompressible Euler.
- A new polar factorization theorem.

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Generalisation of Brenier's proof

Theorem (P2)

Let $(\varphi(t), r(t))$ be a smooth solution to the geodesic equations on the time interval $[t_0, t_1]$. If $(t_1 - t_0)^2 \langle w, \nabla^2 \Psi_{P(t)}(x, r)w \rangle < \pi^2 ||w||^2$ holds for all $t \in [t_0, t_1]$ and $(x, r) \in C(M)$ and $w \in T_{(x,r)}C(M)$, then for every smooth curve $(\varphi_0(t), r_0(t)) \in \operatorname{Aut}_{vol}(C(M))$ satisfying $(\varphi_0(t_i), r_0(t_i)) = (\varphi(t_i), r(t_i))$ for i = 0, 1 and the condition (*), one has

$$\int_{t_0}^{t_1} \|(\dot{\varphi}, \dot{r})\|^2 \, \mathrm{d}t \le \int_{t_0}^{t_1} \|(\dot{\varphi}_0, \dot{r}_0)\|^2 \, \mathrm{d}t \,, \tag{35}$$

with equality if and only if the two paths coincide on $[t_0, t_1]$.

Define $\delta_0 \stackrel{\text{def.}}{=} \min\{r(x, t) : \text{ injectivity radius at } (\varphi(t, x), r(t, x))\}$, then the condition (*) is:

If the sectional curvature of C(M) can assume both signs or if diam(M) ≥ π, there exists δ satisfying 0 < δ < δ₀ such that the curve (φ₀(t), r₀(t)) has to belong to a δ-neighborhood of (φ(t), r(t)), namely

 $d_{\mathcal{C}(M)}\left((\varphi_0(t,x),r_0(t,x)),(\varphi(t,x),r(t,x))\right)\right) \leq \delta$

for all $(x, t) \in M \times [t_0, t_1]$ where $d_{\mathcal{C}(M)}$ is the distance on the cone.

- If C(M) has non positive sectional curvature, then, for every δ as above, there exists a short enough time interval on which the geodesic will be length minimizing.
- **(3)** If $M = S_d(1)$, the result is valid for every path $(\dot{\varphi}_0, \dot{r}_0)$.

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In short:

Gain w.r.t. Ebin and Marsden

- Ebin and Marsden proved that: Smooth solutions are minimizing in a H^{d/2+2+ε} neighborhood.
- We have: Smooth solutions are minimizing in a W^{1,∞} neighborhood.

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In short:

Gain w.r.t. Ebin and Marsden

- Ebin and Marsden proved that: Smooth solutions are minimizing in a H^{d/2+2+ε} neighborhood.
- We have: Smooth solutions are minimizing in a W^{1,∞} neighborhood.

Corollary (P2)

When $M = S_1$, smooth solutions to the Camassa-Holm equation

$$\begin{cases} \partial_t u - \frac{1}{4} \partial_{txx} u + 3 \partial_x u \, u - \frac{1}{2} \partial_{xx} u \, \partial_x u - \frac{1}{4} \partial_{xxx} u \, u = 0\\ \partial_t \varphi(t, x) = u(t, \varphi(t, x)). \end{cases}$$
(36)

are length minimizing for short times.

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Toward the incompressible Euler equation

Why? Unbalanced OT is linked to standard OT on the cone.

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Toward the incompressible Euler equation

Why? Unbalanced OT is linked to standard OT on the cone.

Question

Understand $\text{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+)$ as a subgroup of $\text{Diff}(\mathcal{C}(M))$?

Answer

The cone $\mathcal{C}(M)$ is a trivial principal fibre bundle over M. The automorphism group Aut $(\mathcal{C}(M)) \subset \text{Diff}(\mathcal{C}(M))$ can be identified with $\text{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+)$. One has $(\varphi, \lambda) : (x, r) \mapsto (\varphi(x), \lambda(x)r)$.

Recall that $\psi \in \operatorname{Aut}(\mathcal{C}(M))$ if $\psi \in \operatorname{Diff}(\mathcal{C}(M))$ and $\forall \lambda \in \mathbb{R}^*_+$ one has $\psi(\lambda \cdot (x, r)) = \lambda \cdot \psi(x, r)$ where $\lambda \cdot (x, r) \stackrel{\text{def.}}{=} (x, \lambda r)$.

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CH as an incompressible Euler equation

The geodesic equation on $\text{Diff}(M) \ltimes C^{\infty}(M, \mathbb{R}^*_+)$ can be extended to $\text{Aut}(\mathcal{C}(M))$ as

$$\frac{D}{Dt}(\dot{\varphi},\dot{\lambda}r) = -\nabla\Psi_P \circ (\varphi,\lambda r), \qquad (37)$$

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where $\Psi_P(x, r) \stackrel{\text{\tiny def.}}{=} r^2 P(x)$.

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where $\Psi_P(x, r) \stackrel{\text{\tiny def.}}{=} r^2 P(x)$.

Question

Does there exist a density $\tilde{\mu}$ on the cone such that $inj(Diff(M)) \subset SDiff_{\tilde{\mu}}(\mathcal{C}(M))$? (answer: yes)

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Question

Does there exist a density $\tilde{\mu}$ on the cone such that $inj(Diff(M)) \subset SDiff_{\tilde{\mu}}(\mathcal{C}(M))$? (answer: yes)

Proof.

The measure $\tilde{\mu} \stackrel{\text{\tiny def.}}{=} r^{-3} dr d\mu$ where μ denotes the volume form on M.

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A new geometric picture



Figure – On the left, the picture represents the Riemannian submersion between $\operatorname{Aut}(\mathcal{C}(M))$ and the space of positive densities on M and the fiber above the volume form is $\operatorname{Aut}_{\operatorname{vol}}(\mathcal{C}(M))$. On the right, the picture represents the automorphism group $\operatorname{Aut}(\mathcal{C}(M))$ isometrically embedded in $\operatorname{Diff}(\mathcal{C}(M))$ and the intersection of $\operatorname{Diff}_{\tilde{\nu}}(\mathcal{C}(M))$ and $\operatorname{Aut}(\mathcal{C}(M))$ is equal to $\operatorname{Aut}_{\operatorname{vol}}(\mathcal{C}(M))$.

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Results

Theorem (P2)

Let φ be the flow of a smooth solution to the Camassa-Holm equation then $\Psi(\theta, r) \stackrel{\text{def.}}{=} (\varphi(\theta), \sqrt{\operatorname{Jac}(\varphi(\theta))}r)$ is the flow of a solution to the incompressible Euler equation for the density $\frac{1}{r^4} r \operatorname{d} r \operatorname{d} \theta$.

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Results

Theorem (P2)

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Case where $M = S_1$, $\mathcal{M}(\varphi) = [(\theta, r) \mapsto r \sqrt{\partial_x \varphi(\theta)} e^{i\varphi(\theta)}]$ then the CH equation is

$$\begin{cases} \partial_t u - \frac{1}{4} \partial_{txx} u \, u + 3 \partial_x u \, u - \frac{1}{2} \partial_{xx} u \, \partial_x u - \frac{1}{4} \partial_{xxx} u \, u = 0\\ \partial_t \varphi(t, x) = u(t, \varphi(t, x)) \,. \end{cases}$$
(38)

The Euler equation on the cone, $C(M) = \mathbb{R}^2 \setminus \{0\}$ for the density $\rho = \frac{1}{r^4}$ Leb is

$$\begin{cases} \dot{v} + \nabla_{v} v = -\nabla p, \\ \nabla \cdot (\rho v) = 0. \end{cases}$$
(39)

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where $v(\theta, r) \stackrel{\text{\tiny def.}}{=} (u(\theta), \frac{r}{2}\partial_x u(\theta)).$

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Toward polar factorization

Definition (Admissible measures)

We say that a positive Radon measure ρ on M is admissible (with respect to vol) if for any $x \in M$, there exists $y \in \text{Supp}(\rho)$ such that $d(x, y) < \pi/2$.

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Consequence (Liero, Mielke, Savaré): Existence of a unique optimal potential which takes finite values a.e. between vol and ρ admissible. Recall that $c(x, y) = -\log(\cos^2(d(x, y) \wedge \pi/2))$.

$$\mathsf{WF}^{2}(\rho_{0},\rho_{1}) = \sup_{(z_{0},z_{1})\in C(M)^{2}} \int_{M} 1 - e^{-z_{0}(x)} \,\mathrm{d}\rho_{0}(x) + \int_{M} 1 - e^{-z_{1}(y)} \,\mathrm{d}\rho_{1}(y)$$
(40)

subject to $\forall (x, y) \in M^2$,

$$z_0(x) + z_1(y) \le -\log\left(\cos^2\left(d(x,y) \wedge (\pi/2)\right)\right)$$
. (41)

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Polar factorization

Theorem (Polar factorization, P2)

Let $(\phi, \lambda) \in \overline{\operatorname{Aut}}(\mathcal{C}(M))$ s.t. $\rho_1 = \pi_0 [(\phi, \lambda), \operatorname{vol}]$ is an absolute continuous admissible measure. Then, there exist a unique minimizer, characterized by a *c*-concave function z_0 , between vol and ρ_1 and a unique measure preserving generalized automorphism $(s, \sqrt{\operatorname{Jac}(s)}) \in \overline{\operatorname{Aut}}_{\operatorname{vol}}(\mathcal{C}(M))$ such that vol a.e.

$$(\phi, \lambda) = \exp^{\mathcal{C}(M)} \left(-\frac{1}{2} \nabla p_{z_0}, -p_{z_0} \right) \circ (s, \sqrt{\operatorname{Jac}(s)})$$
(42)

or equivalently

$$(\phi, \lambda) = \left(\varphi, e^{-z_0}\sqrt{1 + \|\nabla z_0\|^2}\right) \cdot (s, \sqrt{\operatorname{Jac}(s)}), \qquad (43)$$

where $p_{z_0} = e^{z_0} - 1$ and

$$\varphi(x) = \exp_x^M \left(-\arctan\left(\frac{1}{2} \|\nabla z_0(x)\|\right) \frac{\nabla z_0(x)}{\|\nabla z_0(x)\|} \right) \,. \tag{44}$$

Moreover $(s, \sqrt{\operatorname{Jac}(s)})$ is the unique $L^2(M, \mathcal{C}(M))$ projection of (ϕ, λ) onto $\overline{\operatorname{Aut}_{vol}}(\mathcal{C}(M))$.

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Polar factorization

Another formulation of the polar factorization:

Corollary (P2)

Denote by $\mathcal{M}es^{1}(\mathcal{C}(M)))^{\mathbb{R}^{+}_{+}}$ the space of mesurable and approximate differentiable functions $f : \mathcal{C}(M) \mapsto \mathbb{R}$ that satisfy $f(x,r) = r^{2}f(x,1)$ for any $r \in \mathbb{R}^{*}_{+}$. Under the hypothesis of the previous theorem, there exists a unique couple $\left((s, \sqrt{\operatorname{Jac}(s)}), \Psi_{P}\right) \in \overline{\operatorname{Aut}}_{\operatorname{vol}} \times \mathcal{M}es^{1}(\mathcal{C}(M)))^{\mathbb{R}^{*}_{+}}$ such that

$$(\phi, \lambda) = \exp^{\mathcal{C}(M)}(-\nabla \Psi_P) \circ (s, \sqrt{\operatorname{Jac}(s)}), \qquad (45)$$

where $\Psi(x, r) = r^2 z_0(x)$.

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Perspectives

- Study the relaxation of geodesics for CH (uniqueness of the pressure, how the angle of the cone affects the results...)
- Develop numerical approaches following Mérigot et al.
- Treat other fluid dynamic equations ?

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Corresponding polar factorization

Figure - CH equation after the "Madelung transform"

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Corresponding decomposition of vector fields

Polar factorization as extension of the Hodge-Helmholtz decomposition:

$$v = w + \nabla p$$
 where div $(v) = 0$. (46)

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In our case,

$$(v(\theta), r\lambda(\theta)) = \left(w(\theta), \frac{r}{2}\operatorname{div}(w(\theta))\right) + \left(\frac{1}{2}\nabla p(\theta), rp(\theta)\right).$$
 (47)

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A word about smoothness: Monge-Ampère equation

The corresponding Monge-Ampère equation can be written as

$$\det \left[-\nabla^2 z(x) + (\nabla^2_{xx} c)(x, \varphi(x)) \right] = \left| \det \left[(\nabla_{x,y} c)(x, \varphi(x)) \right] \right| e^{-2z(x)} \left(1 + \frac{1}{4} \| \nabla z(x) \|^2 \right) \frac{f(x)}{g \circ \varphi(x)},$$

$$(48)$$

where φ is the *c*-exponential of -z:

$$\varphi(x) = \exp_x^M \left(-\arctan\left(\frac{1}{2} \|\nabla z(x)\|\right) \frac{\nabla z(x)}{\|\nabla z(x)\|} \right) \,. \tag{49}$$

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For the cost $c(x,y) = -\log(\cos^2(d(x,y) \wedge \pi/2))$,

- On the plane, there exist $(x, y) \in M^2$ and $(v, w) \in T_x M \times T_y M$, MTW(x, y, v, w) < 0.
- On the sphere of radius r = 1, as well.
- If r small enough, then numerically, $MTW \ge 0$.

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Link with the reflector problem

Consider the sphere of radius 1/2, then $d(x, y) = \frac{1}{2} \arccos(x \cdot y)$:

$$\begin{aligned} -\log(\cos^2(d(x,y))) &= -\log(1 + \cos(2d(x,y))) + \log(2) \\ &= -\log(1 + x \cdot y) + \log(2) \\ &= -2\log(|x+y|) = 2c_r(x,-y) \end{aligned}$$

The cost for the reflector antenna is $c_r(x, y) = -\log(|x - y|)$. Clearly,

 $\operatorname{sgn}(\operatorname{MTW}(c_r(\cdot,\cdot))) = \operatorname{sgn}(\operatorname{MTW}(c_r(\cdot,-\cdot)))$

Therefore, $MTW(-\log(\cos^2(d))) \ge 0$ on the sphere of radius 1/2. (Loeper, Lee and Li).

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