Linear/nonlinear approaches for the approximation of convection-diffusion equations

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Outline of the talk



2 Two-Point Flux Approximation of linear convection-diffusion



3 Hybrid finite volume schemes

Outline of the talk

1 Basics on convection-diffusion equations

2 Two-Point Flux Approximation of linear convection-diffusion



Fokker-Planck equation (with anisotropy)

$$\begin{cases} \partial_t u + \operatorname{div} \mathbf{J} = 0, \quad \mathbf{J} = \mathbf{\Lambda}(-\nabla u - u\nabla\Phi), \text{ in } \Omega \times (0,T) \\ + \text{ Dirichlet on } \Gamma^D \text{ and no-flux on } \Gamma^N, \ \partial\Omega = \Gamma^D \bigcup \Gamma^N \\ u(\cdot,0) = u_0 \ge 0 \end{cases}$$

Examples

• Semiconductor models, corrosion models • $\Lambda = \mathbf{I}$ or $\Lambda = \frac{1}{1+b^2} \begin{pmatrix} 1 & \pm b \\ \mp b & 1 \end{pmatrix}$ (with magnetic field)

 \twoheadrightarrow coupling with a Poisson equation for Φ

- Porous media flow
 - $woheadrightarrow \Lambda$ bounded, symmetric and uniformly elliptic

$$\Rightarrow \Phi = gz$$

Assumptions : $\Phi \in C^1(\overline{\Omega}, \mathbb{R})$, $\int_{\Omega} u_0 > 0$.

Structural properties

$$\begin{cases} \partial_t u + \text{div } \mathbf{J} = 0, \quad \mathbf{J} = \mathbf{\Lambda}(-\nabla u - u\nabla \Phi), \\ u(\cdot, 0) = u_0 \ge 0 \quad + \text{ boundary conditions} \end{cases}$$

- Existence and uniqueness of the solution
- Nonnegativity of u, mass conservation if $\Gamma^D = \emptyset$
- Existence of a thermal equilibrium :

$$u_{\infty} = \rho e^{-\Phi} (\Longrightarrow \mathbf{J} = 0)$$

 $\Rightarrow \text{ if } \Gamma^D = \emptyset, \\ \rho = \frac{\int_{\Omega} u_0}{\int_{\Omega} e^{-\Phi}}, \quad \text{so that} \quad \int_{\Omega} u_\infty = \int_{\Omega} u_0$

 $\implies \text{if } \Gamma^D \neq \emptyset \text{ and } u^D = \rho e^{-\Phi^D} \text{ on } \Gamma^D.$

Reformulation of the convection-diffusion fluxes

$$\begin{cases} \partial_t u + \text{div } \mathbf{J} = 0, \quad \mathbf{J} = -\nabla u - u\nabla\Phi, \\ \mathbf{J} \cdot \mathbf{n} = 0 \text{ on } \Gamma^N \text{ and } u = u^D \text{ on } \Gamma^D. \end{cases} \qquad \qquad u^{\infty} = \rho e^{-\Phi}$$

Equivalence

$$\mathbf{J} = -\nabla u - u\nabla\Phi$$
$$= -u^{\infty}\nabla\frac{u}{u^{\infty}}$$
$$= -u\nabla\log\frac{u}{u^{\infty}}$$
$$= -u\nabla(\log u + \Phi)$$

Towards nonlinear convection diffusion fluxes (à la Onsager)

$$\mathbf{J} = -\eta(u)\nabla(\mu(u) + z_u\Phi)$$

 η : mobility, μ : chemical potential, z_u : charge

Long time behaviour of the Fokker-Planck equation

$$\begin{cases} \partial_t u + \operatorname{div} \mathbf{J} = 0, \quad \mathbf{J} = -\nabla u - u\nabla\Phi, \\ \mathbf{J} \cdot \mathbf{n} = 0 \text{ on } \Gamma^N \text{ and } u = u^D \text{ on } \Gamma^D. \end{cases} \qquad u^{\infty} = \rho e^{-\Phi} \end{cases}$$

Dissipation of some relative $\Psi\text{-entropies}$

 Ψ $C^2\text{-convex}$ function, $\Psi(1)=\Psi'(1)=0$

$$\begin{split} \mathbb{E}(t) &= \int_{\Omega} u^{\infty} \Psi(\frac{u}{u^{\infty}}) \\ \mathbb{I}(t) &= \int_{\Omega} u^{\infty} \nabla \Psi'(\frac{u}{u^{\infty}}) \cdot \nabla \frac{u}{u^{\infty}} \\ &= \int_{\Omega} u \nabla \Psi'(\frac{u}{u^{\infty}}) \cdot \nabla \log \frac{u}{u^{\infty}} \end{split} \right\} \qquad \begin{array}{l} \frac{d\mathbb{E}}{dt} + \mathbb{I} = 0, \\ & \text{with } \mathbb{I} \geq 0. \end{split}$$

Exponential decay of the relative Ψ -entropies

For some specific choice of Ψ and thanks to functional inequalities,

$$\exists \nu \in \mathbb{R}, \quad \mathbb{E}(t) \le \mathbb{E}(0)e^{-\nu t}.$$

Outline of the talk

Basics on convection-diffusion equations

2 Two-Point Flux Approximation of linear convection-diffusion



TPFA finite volume schemes

 $\partial_t u + \operatorname{div} \mathbf{J} = 0$



Generic form of the finite volume schemes

• (forward Euler) scheme in time

$$\frac{u^{n+1}-u^n}{\Delta t} + \operatorname{div}\, \mathbf{J}^{n+1} = 0$$

integration of the balance law over control volumes

$$\begin{cases} \mathsf{m}(K) \frac{u_{K}^{n+1} - u_{K}^{n}}{\Delta t} + \sum_{\sigma \in \mathcal{E}_{K}} \mathcal{F}_{K,\sigma}^{n+1} = 0, \forall K \in \mathcal{T} \\ \mathcal{F}_{K,\sigma} \approx \int_{\sigma} \mathbf{J} \cdot \mathbf{n}_{K,\sigma} \quad \forall \sigma \in \mathcal{E}_{K}. \end{cases}$$

First examples of numerical fluxes (TPFA)

•
$$\mathbf{J} = -\nabla u$$
 $\mathcal{F}_{K,\sigma} = \mathbf{m}(\sigma) \frac{u_K - u_L}{\mathbf{d}_{\sigma}} = \tau_{\sigma}(u_K - u_L)$
• $\mathbf{J} = \mathbf{v}u$ $\mathcal{F}_{K,\sigma} = \mathbf{m}(\sigma) v_{K,\sigma} \frac{u_K + u_L}{2}, \ v_{K,\sigma} \approx \oint_{\sigma} \mathbf{v} \cdot \mathbf{n}_{K,\sigma}$





• upwind flux : $B_{up}(x) = 1 + \max(-x, 0)$ • centered flux : $B_{ce}(x) = 1 - \frac{x}{2}$ • SG flux : $B_{SG}(x) = \frac{x}{\exp x - 1}$

□ C.-H., DRONIOU, 2011



• $S_{1,-1}(x,y) = \sqrt{xy}$ • $S_{0,-1} = xy(\log x - \log y)/(x-y)$ • $S_{2,1} = (x+y)/2$ • $S_{-2,-1} = 2xy/(x+y)$

Heida, Kantner, Stefan, 2021Brezzi, Marini, Pietra, 1989

5-numerical fluxes / exp. fitting schemes

$$\mathcal{F}_{K,\sigma} \approx \int_{\sigma} \mathbf{J} \cdot \mathbf{n}_{K,\sigma}$$
with $\mathbf{J} = -u^{\infty} \nabla \frac{u}{u^{\infty}}, \quad u^{\infty} = e^{-\Phi}.$

$$\mathcal{F}_{K,\sigma} = \tau_{\sigma} S(u_{K}^{\infty}, u_{L}^{\infty}) \left(\frac{u_{K}}{u_{K}^{\infty}} - \frac{u_{L}}{u_{L}^{\infty}}\right),$$
with S Stolarsky mean : $S_{\alpha,\beta}(x, y) = \left(\frac{\beta}{\alpha} \frac{x^{\alpha} - y^{\alpha}}{x^{\beta} - y^{\beta}}\right)^{\frac{1}{\alpha - \beta}}$

Properties of the $S\mbox{-}{\rm functions}$

$$S(x,y) = xS(1,\frac{y}{x})$$
$$S(x,y) = yS(1,\frac{x}{y})$$
$$\Longrightarrow \begin{cases} \frac{S(u_K^{\infty}, u_L^{\infty})}{u_K^{\infty}} = S(1, e^{-(\Phi_L - \Phi_K)}) \\ \frac{S(u_K^{\infty}, u_L^{\infty})}{u_L^{\infty}} = S(1, e^{-(\Phi_K - \Phi_L)}) \end{cases}$$

5-numerical fluxes / exp. fitting schemes

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S-flux vs B-flux

The S-flux rewrites as a B-flux with

$$B(x) = S(1, e^{-x})$$



B-scheme vs S-scheme?

• S-flux \implies B-flux, with $B(x) = S(1, e^{-x})$,

but B(x) - B(-x) = -x iff $S = S_{0,-1}$ and $B = B_{SG}$.

• B-flux
$$\implies$$
 S-flux ?

if $S(u_K^{\infty}, u_L^{\infty}) = u_K^{\infty} B(\Phi_L - \Phi_K) = u_L^{\infty} B(\Phi_K - \Phi_L)$ iff $B = B_{SG}$, and therefore $S = S_{0,-1}$.

- Preservation of the thermal equilibrium ? always true for the S-flux.
- Preservation of the positivity ?
 yes, for both schemes (monotonicity of the fluxes).
- Exponential decay towards thermal equilibrium ? discrete counterpart of the entropy-dissipation estimate ?

Reformulation of the SG B-fluxes

$$\mathbf{J} = -u\nabla\log\frac{u}{u^{\infty}} = -u\nabla(\log u + \Phi)$$

With the Bernoulli function $B(x) = x/(e^x - 1)$

$$\mathcal{F}_{K,\sigma} = \tau_{\sigma} \Big(B \big(\Phi_L - \Phi_K \big) u_K - B \big(-\Phi_L + \Phi_K \big) u_L \Big)$$
$$0 = \tau_{\sigma} \Big(B (\log u_K - \log u_L) u_K - B (\log u_L - \log u_K) u_L \Big)$$

This implies

$$\mathcal{F}_{K,\sigma} = \tau_{\sigma} \left(\frac{B(y) - B(x)}{x - y} u_K + \frac{B(-x) - B(-y)}{x - y} u_L \right) (x - y)$$

with
$$x = \log u_K - \log u_L$$
, $y = \Phi_L - \Phi_K$,
so that $x - y = \log u_K + \Phi_K - (\log u_L + \Phi_L)$

□ CANCÈS, C.-H., FUHRMANN, GAUDEUL, 2021

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$$0 = \tau_{\sigma} \Big(B (\log u_K - \log u_L) u_K - B (\log u_L - \log u_K) u_L \Big)$$

This implies

$$\mathcal{F}_{K,\sigma} = \tau_{\sigma} r_{KL} \Big(\log u_K + \Phi_K - (\log u_L + \Phi_L) \Big)$$

with r_{KL} convex combination of u_K and u_L (depending also on Φ_K and Φ_L)

□ Cancès, C.-H., Fuhrmann, Gaudeul, 2021

More generally : nonlinear numerical fluxes

$$\mathbf{J} = -\nabla u - u\nabla \Phi = -u\nabla(\log u + \Phi)$$
$$\mathcal{F}_{K,\sigma} \approx \int_{\sigma} -u\nabla(\log u + \Phi) \cdot \mathbf{n}_{K,\sigma}$$

$$\mathcal{F}_{K,\sigma} = \tau_{\sigma} r(u_K, u_L) \Big(\log u_K + \Phi_K - \log u_L - \Phi_L \Big)$$

with $r(u_K, u_L)$ a given mean value between u_K and u_L

\Box Cancès, C.-H., Herda, Krell, 2020

Beyond the TPFA schemes

Drawbacks of the TPFA schemes

- Admissibility of the mesh (orthogonality property)
- $\Lambda = I$

Main objectives from now on

- Design of schemes that are applicable
 - on almost-general meshes,
 - for anisotropic equations,
- while preserving :
 - positivity, conservation of mass,
 - thermal equilibrium and long-time behaviour,
- and with the possibility of extension to high order schemes.

→ Hybrid finite volume schemes.

□ Eymard, Gallouët, Herbin, 2010

Outline of the talk



3 Hybrid finite volume schemes

Mesh and unknowns

Mesh $\mathcal{D} = (\mathcal{M}, \mathcal{E}, \mathcal{P})$

• \mathcal{M} : set of the control volumes (K)

- \mathcal{E} : set of the faces (σ)
- \mathcal{P} : set of the cell centers $(x_K)_{K \in \mathcal{M}}$



Set of discrete unknowns $\underline{V}_{\mathcal{D}}$

 $\underline{u}_{\mathcal{D}} = \{(u_K)_{K \in \mathcal{M}}, (u_{\sigma})_{\sigma \in \mathcal{E}}\} \text{ and } u_{\mathcal{M}} : \Omega \to \mathbb{R}$

Mesh and unknowns

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Set of discrete unknowns $V_{\mathcal{D}}$

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HFV scheme for a diffusion equation

 $-{
m div}~({f \Lambda}
abla u) = f ~~+$ boundary conditions

Foundations of the HFV scheme

- A discrete gradient operator ∇_D on <u>V</u>_D :
 * ∇_D<u>v</u>_D is piecewise constant on the pyramidal submesh,
 * on P_{K,σ}, ∇_D<u>v</u>_D depends only on <u>v</u>_K = (v_K, (v_σ)_{σ∈E_K}) and is made of a consistent part and a stabilisation part.
- Some discrete bilinear forms $a_{\mathcal{D}}^{\Lambda}$ and $(a_{K}^{\Lambda})_{K\in\mathcal{M}}$:

$$a_{\mathcal{D}}^{\Lambda}(\underline{u}_{\mathcal{D}}, \underline{v}_{\mathcal{D}}) = (\Lambda \nabla_{\mathcal{D}} \underline{u}_{\mathcal{D}}, \nabla_{\mathcal{D}} \underline{v}_{\mathcal{D}})_{\Omega} = \sum_{K \in \mathcal{M}} a_{K}^{\Lambda}(\underline{u}_{K}, \underline{v}_{K}).$$

→ definition of the scheme via a variational formulation.

HFV scheme for a diffusion equation

Local discrete bilinear forms and numerical fluxes

$$a_{K}^{\Lambda}(\underline{u}_{K}, \underline{v}_{K}) = (v_{K} - v_{\sigma})_{\sigma \in \mathcal{E}_{K}} \cdot \mathbb{A}_{K}(u_{K} - u_{\sigma})_{\sigma \in \mathcal{E}_{K}}$$

with $\mathbb{A}_{K} = (A_{K}^{\sigma, \sigma'})_{\sigma, \sigma' \in \mathcal{E}_{K}} \in \mathcal{S}_{|\mathcal{E}_{K}|}^{++}$
 $F_{K, \sigma}^{\Lambda}(\underline{u}_{\mathcal{D}}) = \sum_{\sigma' \in \mathcal{E}_{K}} A_{K}^{\sigma, \sigma'}(u_{K} - u_{\sigma'})$

The HFV scheme

$$\begin{cases} \sum_{\sigma \in \mathcal{E}_{K}} F_{K,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}}) = \int_{K} f, \\ F_{K,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}}) + F_{L,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}}) = 0 \quad \forall \sigma = K | L, \\ u_{\sigma} = 0 \quad \forall \sigma \in \mathcal{E}_{ext}^{D} \quad (u^{D} = 0), \quad F_{K,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}}) = 0 \quad \forall \sigma \in \mathcal{E}_{ext}^{N}. \end{cases}$$

→ a linear system of equations on $(u_{\sigma})_{\sigma \in \mathcal{E}}$ (after elimination of the cell unknowns) Linear HFV schemes for an advection-diffusion equation

The Hybrid Mixed Method $\mathbf{J} = -\mathbf{\Lambda} \nabla u + \mathbf{w} u$ \Box Beirão da Vega, Droniou, Manzini, 2011

- Keep the same diffusive fluxes $F_{K,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}})$,
- Define some convective fluxes $F_{K,\sigma}^c(\underline{u}_{\mathcal{D}}) \approx \int_{\sigma} u \mathbf{w} \cdot \mathbf{n}_{K,\sigma}$,
- Write the balance law and the conservativity of the fluxes.

The exponential fitting scheme $\mathbf{J} = -\mathbf{\Lambda}(
abla u + u
abla \Phi)$

- Rewrite $\mathbf{J} = \omega \mathbf{\Lambda} \nabla \rho$ with $\rho = \frac{u}{\omega}$, $\omega = e^{-\Phi}$
- Write the HFV scheme for this anisotropic diffusion equation, with a specific averaging of $\omega \Lambda$ on the pyramids $P_{K,\sigma}$.
- Solve the linear system either in the ρ or in the u variable.

Properties of the linear schemes

	НММ	Exponential fitting
Well-posedness	needs coercivity or smallness of the mesh 🖌	~
Preservation of the thermal equilibrium	×	~
Positivity	×	×
Mass conservation	~	~
Asymptotic stability $\ u_{\mathcal{M}}^n - u_{\mathcal{M}}^{\infty}\ _{L^2(\Omega)}$	~	~

Introduction of a nonlinear HFV scheme

div
$$(\mathbf{J}) = f$$
, $\mathbf{J} = -u\mathbf{\Lambda}\nabla(\log u + \Phi)$

Principles of the nonlinear HFV scheme

• Define the nonlinear numerical fluxes :

$$\mathcal{G}_{K,\sigma}^{\Lambda}(\underline{u}_{\mathcal{D}},\underline{\Phi}_{\mathcal{D}}) = r_{K}(\underline{u}_{\mathcal{D}})F_{K,\sigma}^{\Lambda}(\log \underline{u}_{\mathcal{D}} + \underline{\Phi}_{\mathcal{D}})$$

with
$$r_K(\underline{u}_D) = \frac{1}{2} \left(u_K + \frac{1}{|\mathcal{E}_K|} \sum_{\sigma \in \mathcal{E}_K} u_\sigma \right)$$

.

- Write balance law and conservativity of the fluxes.
- C.-H., HERDA, LEMAIRE, MOATTI, 2023
 MOATTI 2023 (PhD thesis)

About the nonlinear scheme

$$\begin{cases} \partial_t u + \operatorname{div} \mathbf{J} = 0, \quad \mathbf{J} = -u\mathbf{\Lambda}\nabla(\log u + \Phi) = -u\mathbf{\Lambda}\nabla\log\frac{u}{\omega}, \\ u(\cdot, 0) = u_0 \ge 0 \quad + \text{ homogeneous Neumann boundary conditions,} \\ M = \int_{\Omega} u_0 > 0 \end{cases}$$

The scheme under its compact form

$$\frac{1}{\Delta t}(u_{\mathcal{M}}^{n}-u_{\mathcal{M}}^{n-1},v_{\mathcal{M}})_{\Omega}+T_{\mathcal{D}}(\underline{u}_{\mathcal{D}}^{n},\underline{w}_{\mathcal{D}}^{n},\underline{v}_{\mathcal{D}})=0\quad\forall\underline{v}_{\mathcal{D}}\in\underline{V}_{\mathcal{D}},\\ \underline{w}_{\mathcal{D}}=\log\frac{\underline{u}_{\mathcal{D}}}{\underline{\omega}_{D}}\\ T_{\mathcal{D}}(\underline{u}_{\mathcal{D}}^{n},\underline{w}_{\mathcal{D}},\underline{v}_{\mathcal{D}})=\sum_{K\in\mathcal{M}}r_{K}(\underline{u}_{K})a_{K}^{\Lambda}(\underline{w}_{K},\underline{v}_{K}).$$

First a priori estimate : conservation of mass

Choose
$$\underline{v}_{\mathcal{D}} = \underline{1}_D$$
 in order to obtain
$$\int_{\Omega} u_{\mathcal{M}}^n = \int_{\Omega} u_{\mathcal{M}}^{n-1} \text{ for all } n \ge 1.$$

About the nonlinear scheme

Second a priori estimate : entropy-dissipation estimate

• Discrete relative entropy :

$$\mathbb{E}^n = \int_{\Omega} u_{\mathcal{M}}^{\infty} \Psi(\frac{u_{\mathcal{M}}^n}{u_{\mathcal{M}}^{\infty}}) \text{ with } \Psi(s) = s \log s - s + 1.$$

• Discrete dissipation :

$$\mathbb{D}^n = T_{\mathcal{D}}(\underline{u}_{\mathcal{D}}^n, \underline{w}_{\mathcal{D}}^n, \underline{w}_{\mathcal{D}}^n).$$

• With the test function $\underline{v}_{\mathcal{D}} = \underline{w}_{\mathcal{D}}^n$ in the scheme, we get

$$\frac{\mathbb{E}^{n+1} - \mathbb{E}^n}{\Delta t} + \mathbb{D}^{n+1} \le 0$$

if the scheme has a positive solution.

Main results

Theorem 1

Existence of a positive discrete solution to the nonlinear scheme.

Theorem 2

Exponential decay of the discrete entropy in time :

$$\mathbb{E}^{n+1} \le (1 + \tilde{\nu} \Delta t)^{-1} \mathbb{E}^n, \forall n \ge 0$$

Exponential decay of the L^1 -distance to the equilibrium

$$\|u_{\mathcal{M}}^n - u_{\mathcal{M}}^{\infty}\|_{L^1(\Omega)} \le C(1 + \tilde{\nu}\Delta t)^{-\frac{n}{2}}.$$

Numerical results

Long-time behaviour (on a fine Kershaw mesh)



Numerical results

Accuracy of stationary solutions (advection-dominated test case)



Concluding remarks

About the non linear hybrid scheme

- Well posedness, positivity and preservation of the equilibrium
- Exponential decay towards the equilibrium
- A first step towards Hybrid High Order schemes
- \Box Lemaire, Moatti, 2024
- Possible extension to nonlinear convection-diffusion equations

TPFA schemes and nonlinear convection-diffusion fluxes

$$\mathbf{J} = -\eta(u)\nabla(\mu(u) + z_u\Phi)$$

- □ Cancès, C.H., Fuhrmann, Gaudeul, 2021
- \Box Cancès, Venel, 2023
- □ Cancès, Herda, Massimini, 2023